



First Space-VLBI Observations of Sagittarius A*

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Received 2021 October 9; revised 2021 November 10; accepted 2021 November 10; published 2021 November 25

Abstract

We report results from the first Earth-space VLBI observations of the Galactic Center supermassive black hole, Sgr A*. These observations used the space telescope Spektr-R of the RadioAstron project together with a global network of 20 ground telescopes, observing at a wavelength of 1.35 cm. Spektr-R provided baselines up to 3.9 times the diameter of the Earth, corresponding to an angular resolution of approximately $55 \mu\text{as}$ and a spatial resolution of $5.5R_{\text{Sch}}$ at the source, where $R_{\text{Sch}} \equiv 2GM/c^2$ is the Schwarzschild radius of Sgr A*. Our short ground baseline measurements ($\lesssim 80 \text{ M}\lambda$) are consistent with an anisotropic Gaussian image, while our intermediate ground baseline measurements (100–250 $\text{M}\lambda$) confirm the presence of persistent image substructure in Sgr A*. Both features are consistent with theoretical expectations for strong scattering in the ionized interstellar medium, which produces Gaussian scatter-broadening on short baselines and refractive substructure on long baselines. We do not detect interferometric fringes on any of the longer ground baselines or on any ground–space baselines. While space-VLBI offers a promising pathway to sharper angular resolution and the measurement of key gravitational signatures in black holes, such as their photon rings, our results demonstrate that space-VLBI studies of Sgr A* will require sensitive observations at submillimeter wavelengths.

Unified Astronomy Thesaurus concepts: [Interstellar scattering \(854\)](#); [Supermassive black holes \(1663\)](#); [Very long baseline interferometry \(1769\)](#)

1. Introduction

Very long baseline interferometry (VLBI) provides the highest direct angular resolution in astronomy, recently culminating in the first image of a black hole (The Event Horizon Telescope Collaboration, et al. 2019). Images of black holes are expected to have a rich interferometric response on long baselines, including small-scale power from both their irregular accretion flows (e.g., Roelofs et al. 2017; Medeiros et al. 2018; Gelles et al. 2021) and from their distinctive gravitationally lensed features such as the “photon ring” (e.g., Gralla 2020; Johnson et al. 2020). A crucial target for VLBI studies is the Milky Way’s nuclear black hole, Sagittarius A* (Sgr A*), which subtends the largest angular size of any known black hole. The intrinsic and scattered structure of Sgr A* have been intensively studied using VLBI at millimeter and centimeter wavelengths (e.g., Jauncey et al. 1989; Alberdi et al. 1993; Lo et al. 1998; Krichbaum et al. 1998; Bower et al. 2004; Shen et al. 2005; Doeleman et al. 2008; Lu et al. 2011; Johnson et al. 2018; Issaoun et al. 2019). A major challenge for these studies is the unusually strong interstellar scattering of Sgr A*.

Kellermann et al. (1977) reported early VLBI observations of Sgr A* at $\lambda = 3.8 \text{ cm}$, which indicated that 25% of the emission came from a compact component, $\lesssim 1 \text{ mas}$ in size. Such a bright, compact component would be detectable using

ground–space VLBI with the orbiting telescope Spektr-R of the RadioAstron program (Kardashev et al. 2013). More recent measurements have not found evidence for this component at centimeter wavelengths (e.g., Lo et al. 1998; Bower et al. 2004; Lu et al. 2011; Bower et al. 2015a; Johnson et al. 2018). However, Gwinn et al. (2014) discovered compact image substructure in Sgr A* at $\lambda = 1.3 \text{ cm}$, with approximately 1% of the total flux density measured on baselines up to 3000 km ($\approx 250 \times 10^6 \lambda$, or 0.8 mas resolution). This signal is consistent with expected “refractive substructure” produced by scattering in the ionized interstellar medium (Narayan & Goodman 1989; Goodman & Narayan 1989; Johnson & Gwinn 2015). Refractive substructure was also detected in AGNs using space-VLBI (Johnson et al. 2016; Pilipenko et al. 2018).

In this Letter, we examine the interferometric properties of Sgr A* on much longer baselines, using ground–space VLBI provided by the orbiting telescope Spektr-R. Launched in 2011, Spektr-R was located on a highly elliptic orbit, with a perigee of 350,000 km and an orbital period of approximately 9 days. It successfully observed using receivers at four wavelengths, extending from 1.3 cm to 92 cm (e.g., Popov et al. 2017; Kovalev et al. 2020). Our observations of Sgr A* include projected baselines 4 times the Earth’s diameter, providing the sharpest view of this source at centimeter wavelengths and the first using space-VLBI.

2. Observations, Correlation, and Processing

We observed Sgr A* on 2015 September 13, using Spektr-R together with a global ground network of 20 telescopes (observing codes RAGS01B and BK193A). To maximize the



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Table 1
Observing Setup and Array for 1.35 cm Observations of Sgr A*

Telescope Name	Telescope Abbreviation	Diameter (m)	Bitrate (Mbps)	IFs per Polarization (MHz)	Observing Time (UT)
Spektr-R	RA	10	128	2 × 16	13.09.2015 01:30—04:30
VLBA St. Croix	SC	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Hancock	HN	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA North Liberty	NL	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Fort Davis	FD	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Los Alamos	LA	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Pie Town	PT	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Kitt Peak	KP	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Owens Valley	OV	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Brewster	BR	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLBA Maunakea	MK	25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
VLA	VLA	27 × 25	2048	4 × 64	12.09.2015 23:00—13.09.2015 04:30
GBT	GBT	100	2048	4 × 64	12.09.2015 23:00—13.09.2015 03:00
LBA ATCA	ATCA	6 × 22	1024	2 × 64	13.09.2015 02:00—12:00
LBA Mopra	MP	22	1024	2 × 64	13.09.2015 02:00—12:00
LBA Ceduna	CD	30	512	1 × 64	13.09.2015 02:00—12:00
LBA Hobart	HO	26	512	1 × 64	13.09.2015 02:00—12:00
KVN Yonsei	KY	21	1024	2 × 64	13.09.2015 07:00—12:00
KVN Ulsan	KU	21	1024	2 × 64	13.09.2015 07:00—12:00
KVN Tamna	KT	21	1024	2 × 64	13.09.2015 07:00—12:00
Tianma	T6	65	512	4 × 32	13.09.2015 07:00—12:00

Note. The VLA and ATCA are linked radio interferometers of 27 and 6 antennas, respectively, which were each phased up and coherently summed for these observations. All telescopes observed dual circular polarizations except Tianmi, which observed only left-hand circular polarization (LCP). All ground telescopes used 2 bit sampling; Spektr-R used 1 bit sampling. Note that the recorded bandwidth and number of intermediate frequency (IF) bands varied significantly across the array, necessitating a multi-stage correlation and analysis procedure.

likelihood of detections on Sgr A*, which has angular broadening from scattering $\propto \lambda^2$, our observations used the highest Spektr-R observing band (a central frequency of 22.2 GHz). In addition, we selected an observing interval where the Earth-space baselines were predominantly north–south, which is the orientation in which the scattering is both weaker and more poorly constrained. Table 1 gives details of the array and observing setup. The time interval of Spektr-R observations was determined mainly by its cooling constraints. The Green Bank Earth station (Ford et al. 2014) was used to collect science data from Spektr-R.

Correlation was performed using a special version of the software correlator `DiFX` (Deller et al. 2011; Bruni et al. 2016) that is capable of dealing with Spektr-R data. Due to scheduling limitations, there was no common scan for all stations on the strong fringe calibrator 1730–130, however, separate scans for the US and LBA-KVN-T6 parts of the array produced fringes for all ground telescopes. Because of the long slew time, Spektr-R did not observe the calibrator. This decision was supported by an extensive experience with RadioAstron operations; onboard telemetry provided comprehensive information on the technical status of the space telescope.

Due to different recorded bandwidths at different antennas, the correlation was performed in three separate passes to recover as much information as possible:

1. native 64 MHz channels for all ground stations excluding T6;
2. 32 MHz channels for all ground stations;
3. 2 × 16 MHz channels for baselines between Spektr-R and GBT, ATCA, and VLA.

Using the VLBI software processing package `PIMA`,⁷ we performed baseline-based fringe fitting, which included determination of the phase acceleration term to statistically separate observations with significant interferometric signal and noise. To do this, we applied the approach described in Petrov et al. (2011) and Kovalev et al. (2020) to estimate the probability of false detection (PFD) for each interferometric fringe obtained. Fringes with $\text{PFD} > 10^{-4}$ were classified as nondetections; for each nondetection, we computed the associated visibility amplitude upper limit.

The ground array data set with 64 MHz channels (IFs) was also processed in `AIPS`⁸ with global fringe fitting to determine single-IF antenna-based solutions. The amplitude was calibrated using a priori gain and system temperature information provided by the observing facilities. We were able to perform and use the global antenna-based fringe fitting with multi-IF solutions for the US portion of the array due to availability and successful application of instrumental phase-calibration data. With 4 IFs, this procedure improved the SNR of our solutions by up to a factor of two.

For our subsequent analysis and discussion, we use results from both these software packages, relying on PFDs and upper limits from `PIMA` for baselines without detected fringes and using the highest possible signal-to-noise ratio of ground antenna-based fringe-fitting solutions in `AIPS`.

Figure 1 shows the baseline coverage of our observations of Sgr A*. This figure also indicates which baselines gave statistically significant interferometric single-baseline fringes using `PIMA`. We did not detect fringes on any baselines to Spektr-R.

⁷ <http://astrogeo.org/pima/>

⁸ <http://www.aips.nrao.edu>

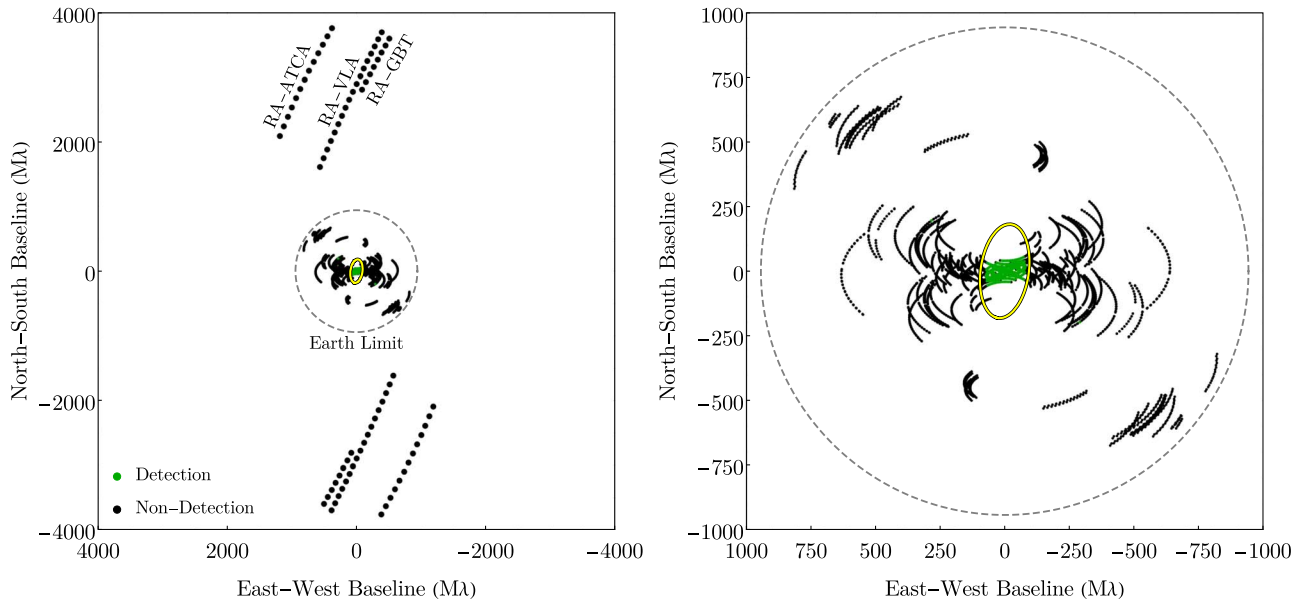


Figure 1. Sgr A* baseline coverage of RadioAstron 1.3 cm observations. The left panel shows the full range of baseline coverage, including only the three most sensitive baselines to Spektr-R; the right panel shows a zoomed-in region with only the ground baselines. Single-baseline detections/non-detections using PIMA are colored green/black. We do not detect Sgr A* on any baselines to Spektr-R. Detections on ground baselines all lie within the region where the scattering kernel is at least 1%, denoted by the yellow ellipse.

3. Analysis and Results

Analysis of our observation was complicated by usually short coherence times and poor phasing of the VLA. To ensure accurate results that recover as much information as possible, we performed our analysis in several stages, using strong detections on short baselines to recover weak detections on long baselines. In all cases, we compare our measurements with predictions from a physically motivated model for the scattering of Sgr A* that was developed by Psaltis et al. (2018). We use source and scattering parameters from Johnson et al. (2018; hereafter, J18), who estimated them by analyzing archival VLBI measurements of Sgr A* extending from 1.3 mm to 29 cm wavelength. The predictions of this model have been successfully tested using VLBI including ALMA at 3.5 mm (Issaoun et al. 2019, 2021) and using monitoring observations with the East Asian VLBI Network at 1.35 and 0.7 cm (I. Cho et al. 2021, in preparation). Specifically, the J18 model has an intrinsic source that is a circular Gaussian with a wavelength-dependent full width at half maximum (FWHM) size $(0.4 \text{ mas}) \times \lambda_{\text{cm}}$, and a scattering screen located 2.7 kpc from the Earth and 5.4 kpc from the source (with a corresponding magnification $M \approx 0.53$). The scattering screen is stochastic, with a shallow scattering power-law index of $\alpha = 1.38$,⁹ an inner scale of $r_{\text{in}} = 800 \text{ km}$, and anisotropic angular broadening at a position angle $\phi_{\text{PA}} = 81^\circ.9$ with FWHMs along the principal axes $\theta_{\text{maj}} = (1.380 \text{ mas}) \times \lambda_{\text{cm}}^2$ and $\theta_{\text{min}} = (0.703 \text{ mas}) \times \lambda_{\text{cm}}^2$ (as $\lambda \rightarrow \infty$).

On baselines that do not significantly resolve Sgr A*, the visibility function is expected to be well approximated by that of the ensemble-averaged image—the product of the visibility function of the intrinsic (unscattered) source and the scattering

kernel (e.g., Coles et al. 1987; Goodman & Narayan 1989). For baselines with $b \lesssim (1 + M)r_{\text{in}} \approx 1200 \text{ km}$, the scattering kernel is well approximated by an elliptical Gaussian (Goodman & Narayan 1989; Psaltis et al. 2018); for baselines that only partially resolve the intrinsic source, the source visibility function is also well approximated by an elliptical Gaussian. Thus, we first fit our short-baseline data using an elliptical Gaussian model described in Section 3.1.

Next, we used the Gaussian fit on short baselines to derive amplitude and phase self-calibration solutions that were applied to all baselines. After discarding data that had no self-calibration solution, we coherently averaged the data in two-minute intervals and across all four IFs to maximize the available sensitivity. We analyzed the averaged data on long ground baselines, which were then sufficiently sensitive to detect a $\lesssim 1\%$ signal, described in Section 3.2.

Finally, we assessed whether the long baseline signal was consistent with the expected refractive noise using the J18 model. Because refractive substructure is partially quenched by an extended source, this provides a test of both the intrinsic and scattering model parameters. However, the comparison is statistical in nature because the signal from refractive noise is stochastic. Our comparisons are described in Section 3.3.

3.1. Short-baseline Analysis: Gaussian Fits

We first fit a Gaussian model to the short-baseline data. To avoid biasing this fit with non-Gaussian contributions from refractive-noise-dominated measurements, we fit only the data on baselines with an expected ensemble-average flux density of at least 2% of the total flux density in the J18 model. The longest baseline meeting this criterion had $|\mathbf{u}| = 97.4 \text{ M}\lambda$ (BR-VLA). We began by coherently averaging the data in frequency (across each IF) and in time (in 2 s intervals). We then applied amplitude and phase self-calibration solutions before coherently averaging the data further, in 2 minute intervals. At this stage, the primary purpose of self-calibration is to avoid phase decoherence in the coherent averaging.

⁹ The power-law index α corresponds to that of the phase-structure function in the inertial range. The 2D power spectrum of phase fluctuations has an index $\beta = \alpha + 2$; the angular broadening scales as $\lambda^{1+2/\alpha}$ at short wavelengths and as λ^2 at long wavelengths. For a review of interstellar scattering, see Rickett (1990) or Narayan (1992).

Table 2
Summary of Gaussian Fits to Sgr A*

λ [cm]	θ_{maj} [μas]	θ_{min} [μas]	P.A. [deg]
1.352	2587 ± 18	1210 ± 136	81.8 ± 0.9
	2580	1394	81.9
1.348	2572 ± 17	1313 ± 115	82.2 ± 0.8
	2566	1387	81.9
1.344	2571 ± 17	1356 ± 112	82.2 ± 0.9
	2552	1381	81.9
1.341	2543 ± 17	1302 ± 116	82.5 ± 0.9
	2538	1374	81.9

Note. Fitted elliptical Gaussian parameters for Sgr A* for each sub-band (for details, see Section 3.1). Below each fitted parameter is its corresponding expected value using the J18 source and scattering parameters.

Next, we fit an elliptical Gaussian model to the averaged visibilities. We fit the model using diagonalized log closure amplitudes (Blackburn et al. 2020), with flat model priors on the major axis, minor axis, and position angle. To avoid non-Gaussian errors from low signal-to-noise ratio (S/N) measurements, we only included closure amplitudes with $S/N > 5$. Because of the wavelength-dependent image size, we fit data from each IF independently. To perform the fit, we utilized `eht-imaging` (Chael et al. 2016a) together with the nested sampling package `dynesty` (Speagle 2020).

Table 2 shows the results of this Gaussian model fitting. Our data tightly constrain all the model parameters, particularly the major axis and position angle. Note that the fitted uncertainties on these parameters only give the uncertainties from thermal noise; refractive scattering causes an additional $\sim 1\%$ variation in the major and minor axis sizes, which is comparable to the error budget from thermal noise for the better-constrained major axis. All our fitted values are consistent with the predictions from the J18 scattering model.

3.2. Long Baseline Analysis: Refractive Noise

After performing the Gaussian fits on short baselines, we derived amplitude and phase self-calibration solutions and applied them to all baselines. For the amplitude self-calibration, we used a total flux density of 1 Jy, which is comparable to the average of historical values at 1.3 cm (e.g., Bower et al. 2015b). This choice of normalization does not affect our remaining analysis, which only relies on the fractional visibility amplitudes relative to the total flux density. In addition, we averaged the data across all IFs to maximize our sensitivity.

Figure 2 shows the resulting visibility measurements, including upper limits on baselines to Spektr-R (computed by PIMA) and highlighting the amplitudes for our most sensitive ground baselines. These baselines are expected to have signals that are dominated by refractive noise; they have S/N up to 8.4 after the final incoherent averaging in time, indicating a reliable detection of image substructure. For each figure, we show the expected envelope of the ensemble-average image and the predicted “renormalized refractive noise” $\hat{\sigma}_{\text{ref}}(\mathbf{u}) = (|\Delta \hat{V}(\mathbf{u})|^2)^{1/2}$, where \hat{V} is the complex visibility

function of the source after centering the image and normalizing the total flux density (for details, see Appendix A of J18).

3.3. Consistency with Expected Properties

We now examine whether the long baseline visibilities are consistent with the expected level of refractive noise for the J18 model. For baselines that resolve the ensemble-average image, refractive noise is well approximated as a circular complex Gaussian random variable (Johnson & Narayan 2016). The visibility amplitude is then drawn from a Rayleigh distribution. The middle 95% of samples drawn from a Rayleigh distribution fall between 0.18 and 2.2 times the mean, $\langle |\Delta \hat{V}(\mathbf{u})| \rangle = \frac{\sqrt{\pi}}{2} \hat{\sigma}_{\text{ref}}(\mathbf{u}) \approx 0.89 \hat{\sigma}_{\text{ref}}(\mathbf{u})$. All the long baseline measurements in Figure 2 are, thus, consistent with the expected level of refractive noise.

We can make more precise comparisons by averaging measurements from different baselines, which sample partially independent elements of the refractive noise. As a simple comparison, we took an incoherent average of all (debiased) visibility amplitudes on baselines longer than 150×10^6 and with thermal noise less than 10 mJy. This procedure gave an average amplitude of 5.7 mJy. We then generated a set of 1000 scattered images using the `stochastic-optics` module of `eht-imaging` (Johnson 2016); we created synthetic observations for each, with baseline coverage and thermal noise matching our observations. For each synthetic observation, we computed the incoherent average on long baselines exactly as was done for the real observation. This procedure gave a median amplitude of 5.4 mJy, with the inner 95% of samples falling between 4.3 and 8.0 mJy. Hence, the average of our long baseline measurements is consistent with the predicted range of refractive noise.

Our nondetections to Spektr-R are likewise consistent with theoretical expectations for a power-law theory of the scattering and substructure (see Figure 2). We do not find evidence for any long baseline visibilities above the expected refractive noise floor, $\sigma_{\text{ref}} \sim 1 \text{ mJy} \times (|\mathbf{u}|/10^9)^{-0.69}$. In particular, this refractive noise is ~ 2 orders of magnitude below the most stringent upper limits of ~ 150 mJy on baselines to Spektr-R.

An unexpected feature of our long ground baseline measurements is that they are quite close in amplitude to those of Gwinn et al. (2014; see Figure 6 of J18), which were taken on 2014 March 7. Namely, both observations measure visibility amplitudes of ~ 4 mJy on baselines from 150–200 M λ , with a rise to ~ 8 mJy on the ~ 250 M λ baselines GBT-BR and GBT-OV. These observations are separated by 554 days, while the expected correlation timescale of the refractive noise is a few months (for a characteristic scattering velocity of 50 km s $^{-1}$). Thus, we expect that the similarity is simply a chance alignment; if future measurements show that this signal is persistent, it would either indicate a much slower scattering velocity or that the measurements are sensitive to the ensemble-average properties of the scattering rather than of refractive substructure.

4. Summary

Space-VLBI is an exciting frontier for black hole astrophysics with the potential to resolve the gravitationally lensed “photon rings” of nearby supermassive black holes (Johnson et al. 2020) to measure the masses of *thousands* of supermassive black holes via their “shadow” diameters (Pesce et al.

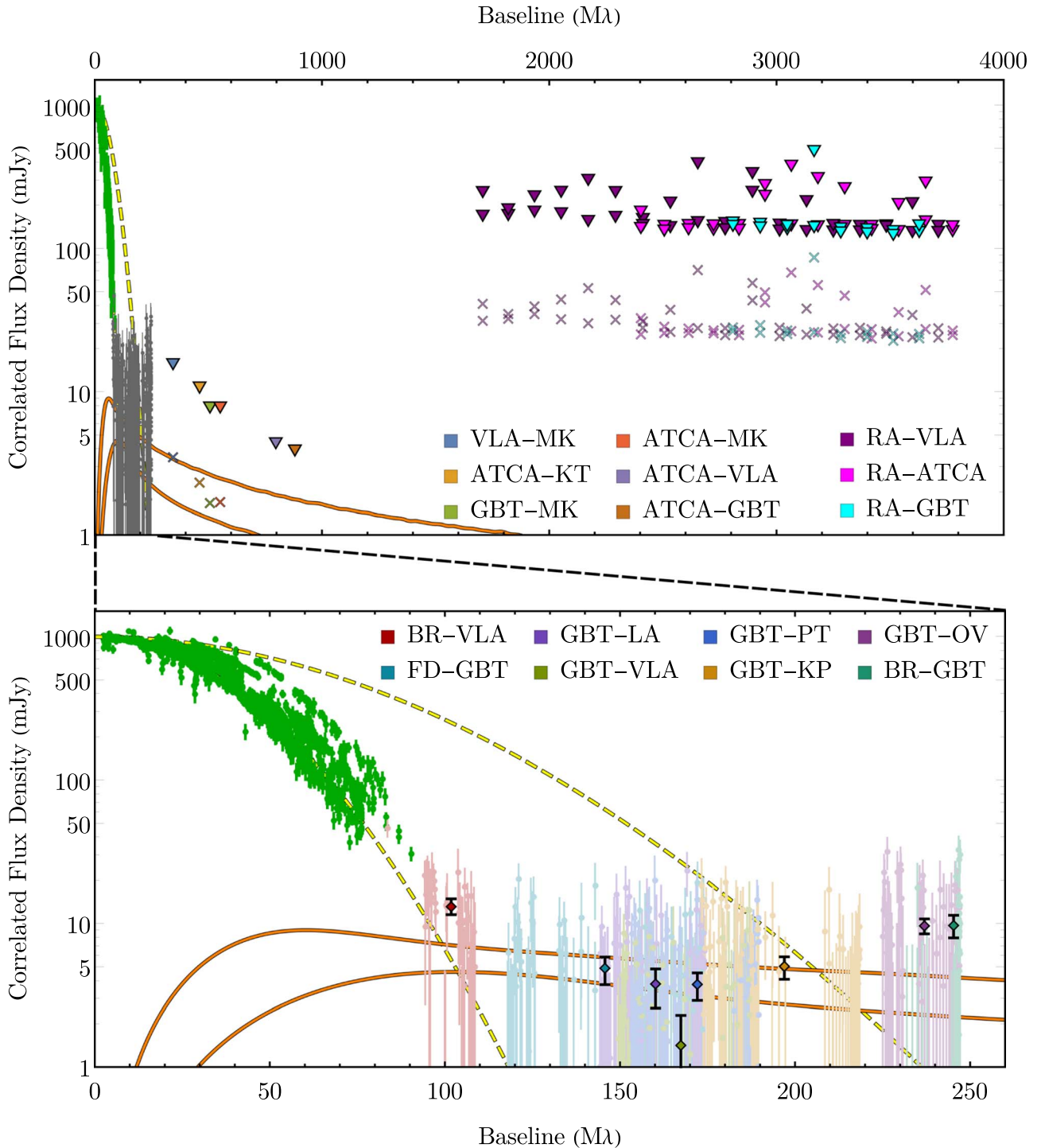


Figure 2. Top: visibility amplitude as a function of baseline length for Sgr A*. Green points show strong detections used to derive the Gaussian self-calibration solution. Gray points show the long baseline visibilities after coherently averaging using the short-baseline self-calibration; for clarity, only visibilities with thermal noise less than 10 mJy and $|u| > 80 \times 10^6$ are shown. Yellow dashed lines show the expected major (lower) and minor (upper) axes of the ensemble-averaged image; orange curves show the expected root-mean-square (renormalized) refractive noise along the major (upper) and minor (lower) axes. Colored triangles show the PIMA fringe amplitudes for representative long baselines, including the three most sensitive ground-space baselines; the corresponding thermal noise for each is indicated with a cross. None of these fringe amplitudes is a statistically significant detection. Bottom: zoom-in on ground baselines. Colored points show a subset of the most sensitive long ground baselines; the incoherent average on each baseline is also shown and indicated with a diamond. Uncertainties shown are $\pm 1\sigma$. Both NL-VLA and BR-VLA have significant flux expected in the ensemble-averaged image, but the remaining long baselines are expected to be dominated by refractive noise.

2021) and to track the orbits of many supermassive black hole binaries (D’Orazio & Loeb 2018). Proposed mission concepts to observe Sgr A* have focused on exploring configurations that allow for rapid baseline sampling to reconstruct movies (e.g., Johnson et al. 2019; Palumbo et al. 2019) as well as

enabling observations at submillimeter wavelengths for which interstellar scattering is sharply reduced (e.g., Roelofs et al. 2019; Kudriashov et al. 2021).

Space-VLBI at longer wavelengths provides crucial input for these designs and gives firm system requirements. In this

Letter, we have presented the first space-VLBI of Sgr A*, observing at $\lambda = 1.35$ cm using Spektr-R and 20 ground antennas. Our short ground baselines are well fit by an elliptical Gaussian image, with parameters matching those of historical measurements. Our long ground baselines confirm the presence of persistent small-scale structure in the scattered image of Sgr A*, originally discovered by Gwinn et al. (2014), at a level that is also consistent with predictions for refractive interstellar scattering. We do not detect interferometric fringes on any baselines to Spektr-R.

Our work highlights the severe challenges for observing Sgr A* with space-VLBI. Because of diffractive interstellar scattering, improving the angular resolution achievable from the ground will likely require observations with $\lambda \lesssim 2$ mm, while detecting features such as the lensed photon ring will require $\lambda \lesssim 0.8$ mm.

We thank Eduardo Ros for useful comments on the manuscript, the anonymous ApJL referee for comments that clarified the manuscript, Amy Mioduszewski for consultations regarding VLA observations in the VLBI mode, and Rebecca Azuly for correlator support at MPIfR. The RadioAstron project was led by the Astro Space Center of the Lebedev Physical Institute of the Russian Academy of Sciences and the Lavochkin Scientific and Production Association under a contract with the State Space Corporation ROSCOSMOS, in collaboration with partner organizations in Russia and other countries. The National Radio Astronomy Observatory and the Green Bank Observatory are facilities of the National Science Foundation operated under cooperative agreement by Associated Universities, Inc. The Australian Long Baseline Array is part of the Australia Telescope National Facility which is funded by the Australian Government for operation as a National Facility managed by CSIRO. This research is based on observations correlated at the Bonn Correlator, jointly operated by the Max Planck Institute for Radio Astronomy (MPIfR), and the Federal Agency for Cartography and Geodesy (BKG). M.D.J. thanks the National Science Foundation (AST-1716536, AST-1935980) and the Gordon and Betty Moore Foundation (GBMF-5278) for financial support of this work. Y.Y.K. is supported in the framework of the State project “Science” by the Ministry of Science and Higher Education of the Russian Federation under the contract 075-15-2020-778.

Facilities: RadioAstron Space Radio Telescope (Spektr-R), VLBA, GBT, VLA, LBA, KVN.


Software: AIPS (Greisen 2003), eht-imaging (Chael et al. 2016b), PIMA (Petrov et al. 2011).

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